# NLO corrections in the initial state parton shower Monte Carlo


**S. JADACH**

**A. Kusina, W. Płaczek, M. Skrzypek and M. Sławińska**

Theory Division, Institute of Nuclear Physics PAN, Kraków, Poland



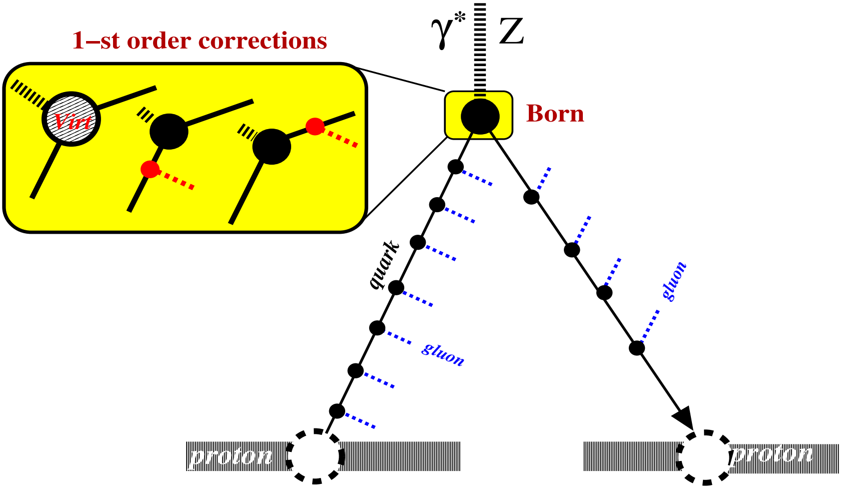
Partly supported by the grants of *Narodowe Centrum Nauki*
DEC-2011/03/B/ST2/02632 and UMO-2012/04/M/ST2/00240
and FP7 grant PITN-GA-2010-264564 (LHCPhenoNet),


Presented at
LHCPhenoNet Annual Meeting, CERN December 2-5th, 2013

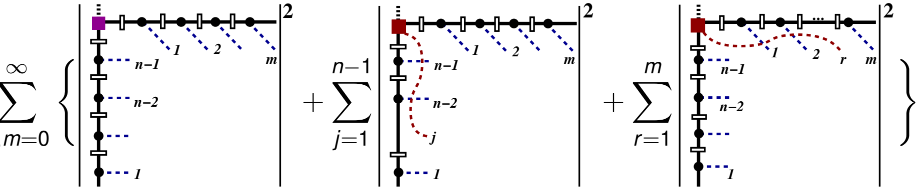
# Outline

- NLO corrections to hard process
  (an alternative to MCatNLO and/or POWHEG)
- NLO corrections in the ladder
  (for NLO parton shower MC + NNLO hard process)

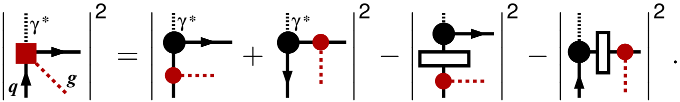
# NLO correcting HARD process



# The essence of our NLO methodology:



NLO real/virtual distributions (subtracted) from Feynman diags



free of any double/single collinear/soft singularities.

# MC weight with NLO corrs. to DY hard proc.

**NLO correction introduced using simple positive MC weight:**

$$W_{MC}^{NLO} = 1 + \Delta_{S+V} + \sum_{j\in F} \frac{\tilde{\beta}_1(\hat{s},\hat{p}_F,\hat{p}_B;a_j,z_{Fj})}{\bar{P}(z_{Fj})\; d\sigma_B(\hat{s},\hat{\theta})/d\Omega} + \sum_{j\in B} \frac{\tilde{\beta}_1(\hat{s},\hat{p}_F,\hat{p}_B;a_j,z_{Bj})}{\bar{P}(z_{Bj})\; d\sigma_B(\hat{s},\hat{\theta})/d\Omega},$$

$\bar{P}(z) \equiv \frac{1+z^2}{2}$. **The IR/Col.-finite real emission part is**

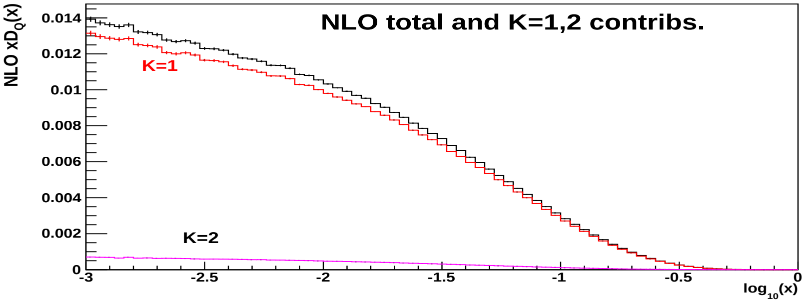
$$\tilde{\beta}_1(\hat{p}_F,\hat{p}_B;q_1,q_2,k) = \left[\frac{(1-\alpha)^2}{2}\frac{d\sigma_B}{d\Omega_q}(\hat{s},\theta_{F1}) + \frac{(1-\beta)^2}{2}\frac{d\sigma_B}{d\Omega_q}(\hat{s},\theta_{B2})\right]$$
$$-\theta_{\alpha>\beta}\frac{1+(1-\alpha-\beta)^2}{2}\frac{d\sigma_B}{d\Omega_q}(\hat{s},\hat{\theta}) - \theta_{\alpha<\beta}\frac{1+(1-\alpha-\beta)^2}{2}\frac{d\sigma_B}{d\Omega_q}(\hat{s},\hat{\theta}),$$

**the kinematics independent virtual+soft correction is**

$$\Delta_{V+S} = \frac{C_F\alpha_s}{\pi}\left(\frac{1}{3}\pi^2 - 4\right) + \frac{C_F\alpha_s}{\pi}\,\frac{1}{2}$$

**Terms like** $\left(\frac{f(z)}{1-z}\right)_+$ **in virt. corrs completely absent!**
**Thanks to transforming MS-bar PDFs to MC Fact scheme,**
**(see also Martin+Ryskin 2013)**

# Only one single term dominates sum in $W_{MC}^{NLO}$

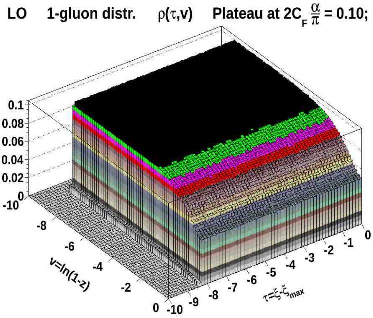



The (-)NLO contributions from $K = 1, 2$ single gluons,
the one with maximum and another one with next-to-max. kT,
in the $x$-distribution of quark entering W boson.

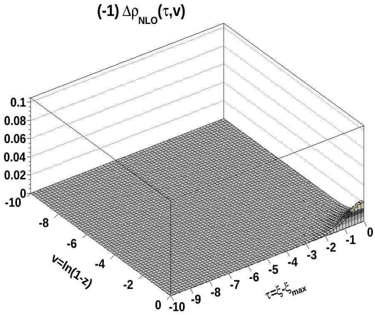
**POWHEG exploits the above. We do it differently, without vetoed/truncated gluons, see next slides...**

# Location and size of the (real) NLO correction on the Sudakov plane (rapidity, ln(1-z))

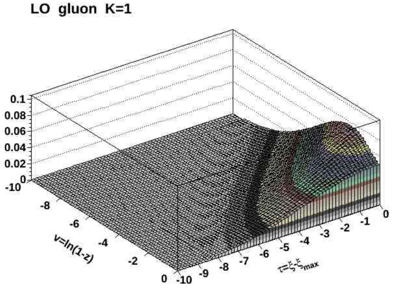
LO inclusive

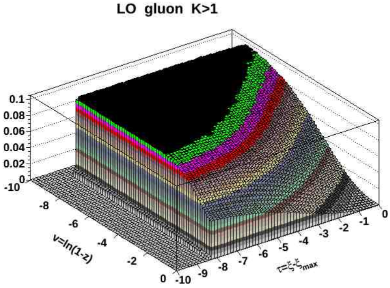



pure NLO



# Gluons generated in angular ordering and re-ordered in kT




Sudakov suppression for the highest kT gluon (K=1)!

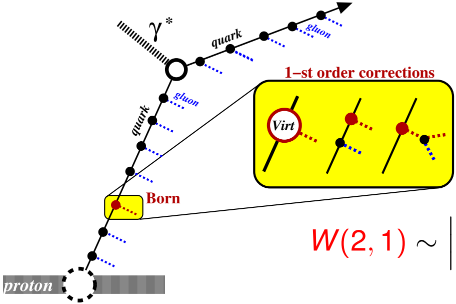
**Our NLO weight with summation is ignorant about the above kT (re-)ordering.**

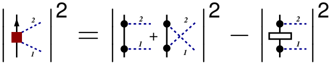
**The summation (or max. kT-selection) takes care of picking up correctly the hardest gluon, for NLO-completeness.**

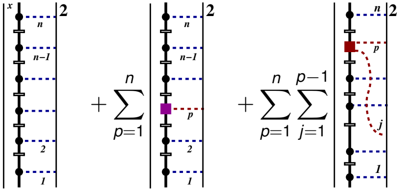
**No need of truncated/vetoed gluon showers.**

# Exclusive NLO-corrections to the middle-of-the-ladder kernels for NLO Monte Carlo parton shower

# NLO-corrected middle-of-the-ladder kernel, $C_F^2$

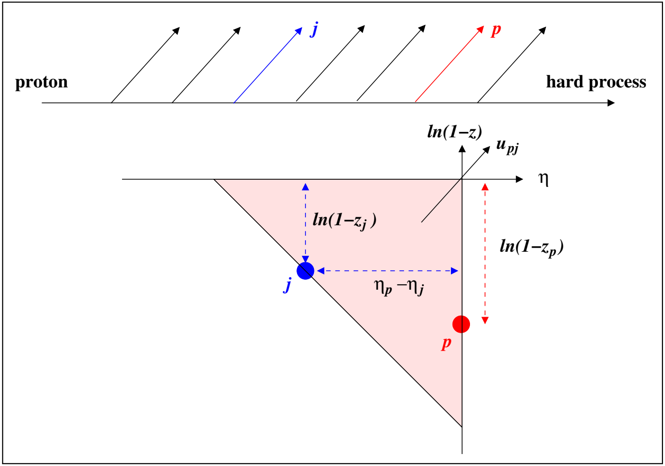



$W(2,1) \sim$ 

$\bar{D}_B^{[1]}(x,Q) = e^{-S_{ISR}} \sum_{n=0}^{\infty} \Big\{$  $\Big\} = e^{-S_{ISR}} \Big\{ \delta_{x=1} +$

$$+ \sum_{n=1}^{\infty} \left( \prod_{i=1}^{n} \int_{Q>a_i>a_{i-1}} d^3\eta_i \; \rho_{1B}^{(1)}(k_i) \right) \left[ 1 + \sum_{p=1}^{n} \beta_0^{(1)}(z_p) + \sum_{p=1}^{n} \sum_{j=1}^{p-1} W(\tilde{k}_p, \tilde{k}_j) \right] \delta_{x=\prod_{j=1}^{n} x_j} \Big\}.$$

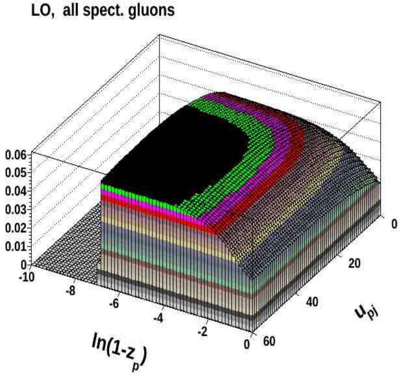
# Define variable $\mathbf{u_{pj}}$ for "u-ordering" in the middle of the ladder

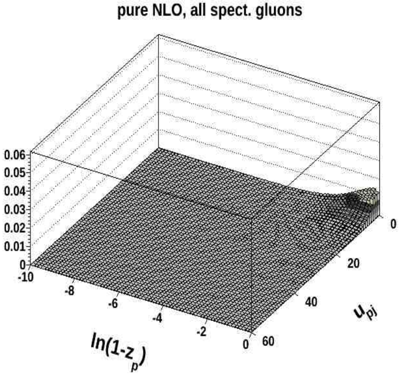



$u_{pj} = |\eta_p - \eta_j| + \lambda \ln(1 - z_j), \quad \lambda \sim 1 - 2.$

Variable $\eta$ is rapitity, $z$ is conventional lightcone variable.

# Location and size of the (real) NLO correction in the ladder on the Sudakov log space

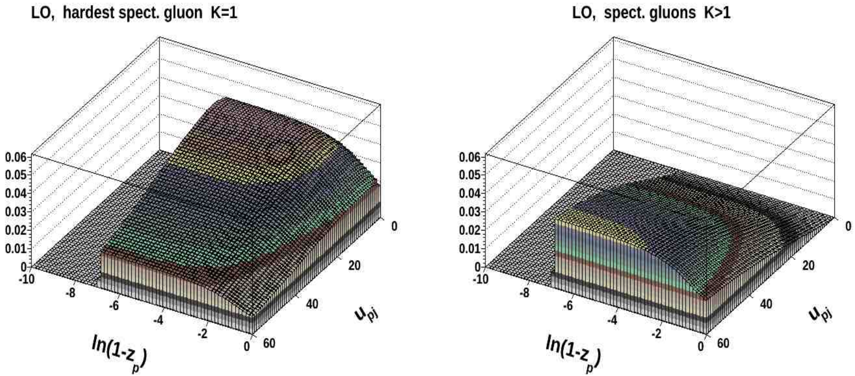



LO inclusive distribution features triple-log IR/coll. singularity, seen as a plateau in 2-dim. projection.

NLO correction IR/coll. finite, nonzero in the corner of the size $\sim 1$.

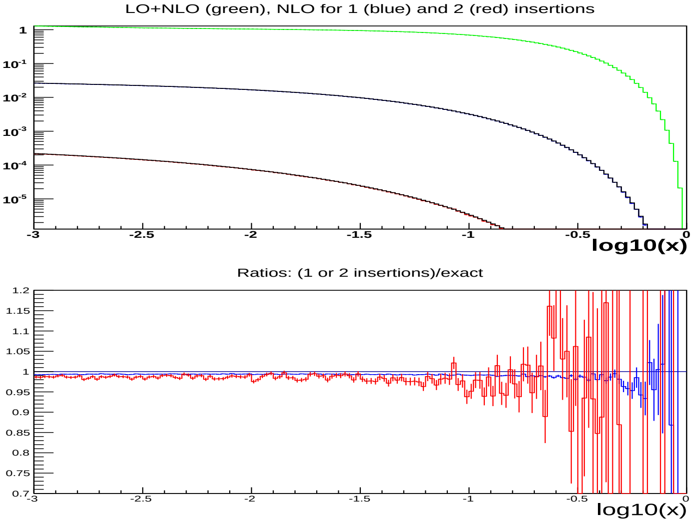
# Split of inclusive LO distribution of gluons into contr. from the u-hardest one and the rest




Distribution of the hardest (in $u_{pj}$) LO spectator gluon approximates well the total distribution in small corner where NLO is non-zero.

# Repetition of test for NLO-corrected ladder

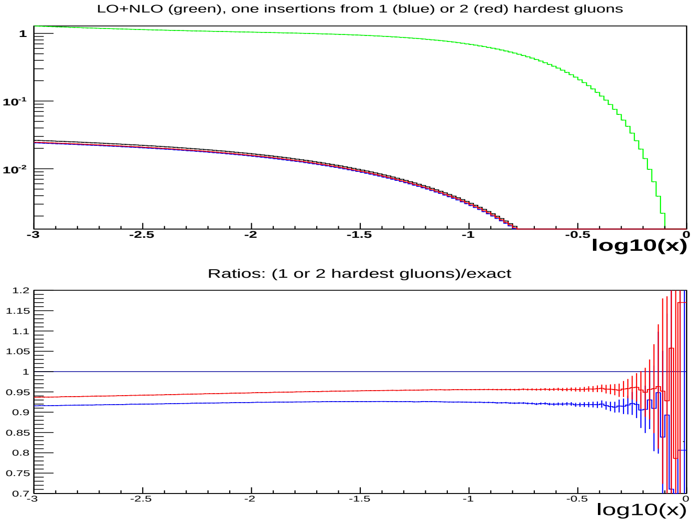
**OLD:** NLO MC vs. analyt. NLO kernels. Perfect agreement




Single ladder, 1GeV-1TeV, 1 or 2 kernels NLO-corrected. (Slow in CPU time).

# Repetition of test for NLO-corrected ladder

**NEW:** NLO contrib. to 1 kernel, 1 and 2 gluons with max. kT




This difference $\sim$ 15% is formally the NNLO/NLO class. (Faster in CPU time).

# Summary

- **An alternative (to MC@NLO or POWHEG) scenario for NLO-corrected hard proc. and LO PSMC is proposed.**
- **Parton shower MC implementing complete NLO DGLAP in the ladders in exclusive way is feasible.**
- Long term: NLO ladder + NNLO hard process,
  (but LO ladder + NLO hard proc. to be optimized first!)
- Most likely application: high quality QCD+EW+QED MC with hard process like $W/Z/H$ boson production.
- Potential gains from new QCD methods are:
  – reducing uncertainties due to PDFs
  – easier implementation of NLO and NNLO corrections to hard process.
  Due to better treatment of “trivial” (albeit numerically sizeable) soft gluon corrections,
  – better environment for low $x$ resumm. (BFKL, CCFM),
  – and more...