

On regularizing the infrared singularities in QCD NLO splitting functions with the new Principal Value prescription

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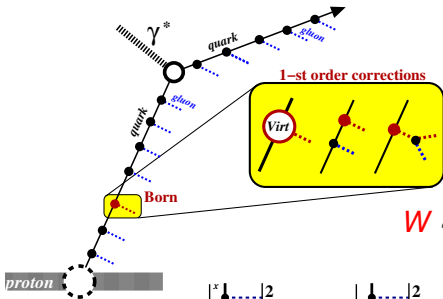


- ▶ Introduction
- ▶ New Principal Value prescription
- ▶ Results for NLO splitting functions
- ▶ Summary

NLO-corrected middle-of-the-ladder kernel, C_F^2



from S. Jadach talk



$$W \sim \left| \begin{array}{c} 2 \\ \text{■} \\ 1 \end{array} \right|^2 = \left| \begin{array}{c} 2 \\ \text{■} + \text{■} \\ 1 \end{array} \right|^2 - \left| \begin{array}{c} 2 \\ \text{■} \\ 1 \end{array} \right|^2$$

$$\bar{D}_B^{[1]}(x, Q) = e^{-S_{ISR}} \sum_{n=0}^{\infty} \left\{ \begin{array}{c} x \\ \bullet \\ n \\ \bullet \\ n-1 \\ \vdots \\ \bullet \\ 2 \\ \bullet \\ 1 \end{array} \right|_2 + \sum_{\rho=1}^n \begin{array}{c} \bullet \\ n \\ \bullet \\ n-1 \\ \vdots \\ \bullet \\ \text{■} \\ p \\ \bullet \\ 2 \\ \bullet \\ 1 \end{array} \right|_2 + \sum_{\rho=1}^n \sum_{j=1}^{\rho-1} \begin{array}{c} \bullet \\ n \\ \bullet \\ n-1 \\ \vdots \\ \bullet \\ \text{■} \\ p \\ \bullet \\ \text{---} \\ j \\ \bullet \\ 2 \\ \bullet \\ 1 \end{array} \right|_2 \Bigg\} = e^{-S_{ISR}} \left\{ \delta_{x=1} + \right.$$

$$\left. + \sum_{n=1}^{\infty} \left(\prod_{i=1}^n \int_{Q > a_i > a_{i-1}} d^3 \eta_i \rho_{1B}^{(1)}(k_i) \right) \left[1 + \sum_{\rho=1}^n \beta_0^{(1)}(z_\rho) + \sum_{\rho=1}^n \sum_{j=1}^{\rho-1} W(\tilde{k}_\rho, \tilde{k}_j) \right] \delta_{x=\prod_{j=1}^n x_j} \right\}.$$

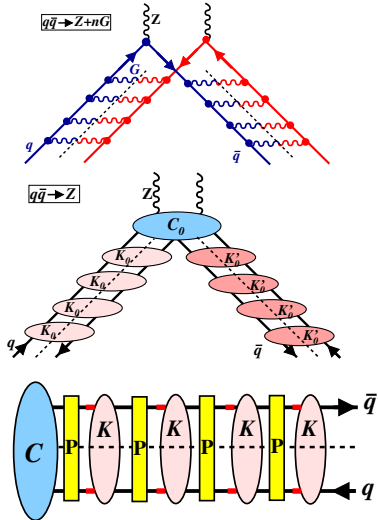
Construction of the evolution kernel in collinear factorization

LO cascade and
construction of the ladder

Include NLO and
group graphs in the ladder into
"two-particle-irreducible" sets K

Use "projection operators" P
to split the ladder and extract kernels

$$\Gamma_{qq} = \text{Tr} \left[\frac{\hat{n}}{4nq} K \hat{p} \right]$$





PV prescription

- ☺ Axial gauge = physical interpretation as parton shower
- ☹ Axial gauge = **spurious (unphysical) singularities**

$$\text{gluon propagator: } \frac{1}{l^2} \left(g^{\mu\nu} - \frac{l^\mu n^\nu + n^\mu l^\nu}{nl} \right)$$

Spurious singularities cancel in full set of diags, but need regularization

Curci, Furmanski, Petronzio [80] Ellis, Vogelsang [96], Heinrich, Kunszt [97]:

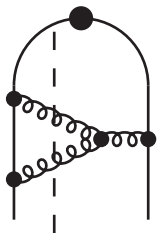
$$\text{Principal Value: } \left[\frac{1}{nl} \right]_{PV} = \frac{nl}{(nl)^2 + \delta^2(pl)^2}$$

PV is more like "phenomenological rule"

Rigorous prescription: Mandelstam [83], Leibbrandt [84].

Difficult in calculations: Bassetto, Heinrich, Kunszt, Vogelsang [97].

Problem with real emission graph



Standard Heinrich, Kunszt [1998]:

$$N(\epsilon, Q^2) \left[\frac{P_{qq}(x)}{\epsilon^3} - 2l_0 \frac{P_{qq}(x)}{\epsilon^2} + \frac{p_{qq}(x)}{\epsilon} \left(-2l_1 + 4l_0 + 2l_0 \ln x - 2l_0 \ln(1-x) \right) \right] + \mathcal{O}\left(\frac{1}{\epsilon}\right)$$

Not good for Parton Shower, needs Real-Virt. cancel.

Parton Shower oriented Jadach, et.al. [2011]:

$$\frac{p_{qq}(x)}{\epsilon} \left(+2l_1 + 4l_0 + 2l_0 \ln x - 2l_0 \ln(1-x) \right) + \mathcal{O}\left(\frac{1}{\epsilon}\right)$$

Good for PS, δ is a cut-off in 4-dimensions, easy to generate in MC:

$$l_0 = \int_0^1 \frac{dx}{[x]_{PV}} \sim \int_\delta^1 \frac{dx}{x} = -\ln \delta, \quad l_1 = \int_0^1 \frac{dx \ln x}{[x]_{PV}} \sim -\frac{1}{2} \ln^2 \delta,$$

$$P_{qq}(x) = p_{qq} + \epsilon(1-x), \quad p_{qq} = \frac{1+x^2}{1-x}$$



New use of PV prescription

Standard: regularize with PV only the gluon propagator

leave other singularities in (+)-component of integration momenta

$$\frac{d^m l}{l_+^{1-\epsilon}}, \quad l_+ = \frac{nl}{np}$$

New proposal: regularize with PV all singularities of the integrand in (+)-component of integration momenta, real & virtual

$$\frac{d^m l}{l_+^{1-\epsilon}} \rightarrow d^m l \left[\frac{1}{l_+} \right]_{PV} \left(1 + \epsilon \ln l_+ + \epsilon^2 \frac{1}{2} \ln^2 l_+ + \dots \right)$$

All (+)-singularities cancel in the final expression (kernel), so extension of "phenomenological PV rule" of Curci-Furmanski-Petronzio possible



Example: virtual three point integral

Must perform (+)-integral as the last one, Ellis, Vogelsang [1996]:

$$\begin{aligned} & \int \frac{d^m l}{(2\pi)^m} \frac{f(l_+)}{l^2(l-q)^2(l-p)^2} = \\ & = \frac{-i}{16\pi^2 q^2} \left(\frac{4\pi}{-q^2} \right)^{-\epsilon} \frac{\Gamma(1-\epsilon)}{-\epsilon} \left[\int_0^x dy f(l_+) \frac{z^\epsilon(1-z)^\epsilon}{1-y} \left(1 + 2\epsilon \ln \frac{1-y}{1-z} \right) \right. \\ & \quad \left. + 2 \frac{\Gamma^2(1+\epsilon)}{\Gamma(1+2\epsilon)} (1-x)^{-\epsilon} \int_x^1 dy f(l_+) (1-y)^{-1+2\epsilon} \right], \end{aligned}$$

$$x = q_+/p_+, \quad y = l_+/p_+, \quad z = y/x, \quad p^2 = (p-q)^2 = 0, \quad m = 4 + 2\epsilon,$$

Singularities at $y = 0$ and $y = x$: only from gluon propagator.

Singularity at $y = 1$: not from gluon propagator!

Proposal: treat all (+)-singularities on equal footing

Note: (+)-singularities lead to $1/\epsilon^3$ poles in kernel

Example: scalar non-axial integral

kinematics: $p^2 = (p - q)^2 = 0$

$$J_3^F = \int \frac{d^m l}{(2\pi)^m} \frac{1}{l^2 (q - l)^2 (p - l)^2}$$

The PV prescription:

$$J_3^F = C \left(-\frac{1}{\epsilon^2} + \frac{\pi^2}{6} \right), \quad C = i \frac{\Gamma(1 - \epsilon)}{(4\pi)^2 |q^2|} \left(\frac{4\pi}{|q^2|} \right)^{-\epsilon}$$

New PV prescription:

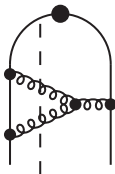
$$J_3^F = C \left(-\frac{2l_0 + \ln(1 - x)}{\epsilon} - 4l_1 + 2l_0 \ln(1 - x) + \frac{\ln^2(1 - x)}{2} \right),$$

$\frac{1}{\epsilon^2}$ replaced by l_1 and $\frac{1}{\epsilon} l_0$

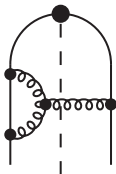
NLO kernels P_{qq} and P_{gg} in New PV scheme



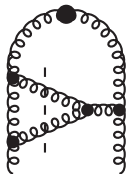
There are only four graphs with $1/\epsilon^3$ singularity:



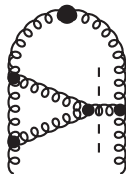
$$\tilde{\Gamma}_{qq}^{(d_R)}(x, \epsilon)$$



$$\tilde{\Gamma}_{qq}^{(d_V)}(x, \epsilon)$$



$$\tilde{\Gamma}_{gg}^{(d_R)}(x, \epsilon)$$



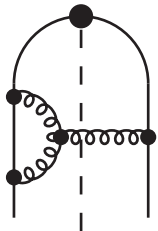
$$\tilde{\Gamma}_{gg}^{(d_V)}(x, \epsilon)$$

related to P_{qq} and P_{gg} splitting functions in a standard way:

$$\tilde{\Gamma}_{qq(gg)}(x, \epsilon) = \delta_{1-x} + \frac{1}{\epsilon} \left(\frac{\alpha_S}{2\pi} P_{qq(gg)}^{LO}(x) + \frac{1}{2} \left(\frac{\alpha_S}{2\pi} \right)^2 P_{qq(gg)}^{NLO}(x) + \dots \right) + \mathcal{O}\left(\frac{1}{\epsilon^2}\right)$$

Virtual contribution to NLO P_{qq} kernel in NPV

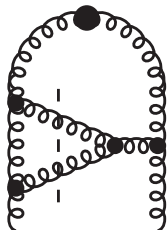
The virtual graph contributes:



$$\begin{aligned} \tilde{\Gamma}_{qq}^{(d_V)}(x, \epsilon) = & -\frac{1}{\epsilon^2} P_{qq}(1 + \epsilon \ln(1-x)) \tilde{Z}_{d_V} \\ & - \frac{1}{\epsilon} p_{qq} \left[l_0(2 \ln x + 2 \ln(1-x)) - 6l_1 - \text{Li}_2(1-x) \right. \\ & \quad \left. + \ln^2 x - 3 + \frac{8}{12} \pi^2 \right] + \frac{1}{\epsilon} \frac{1}{2} \frac{1+x}{1-x}, \\ \tilde{Z}_{d_V} = & 4l_0 + 2 \ln(1-x) + \ln x - \frac{3}{2}, \end{aligned}$$

Inclusive sum of Real and Virtual graphs $\tilde{\Gamma}_{qq}^{(d)}$ identical as in standard PV scheme

Real contribution to NLO P_{gg} kernel in NPV



Only ϵ^{-1} poles

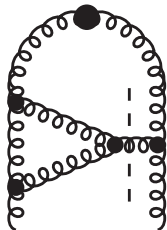
- the calculation can be done in 4-dimensions,
- much simpler than in standard PV

The real graph in New PV prescription:

$$\begin{aligned} \tilde{\Gamma}_{gg}^{(d_R)}(x, \epsilon) = C_S^{(d_R)} \frac{1}{\epsilon} & \left[P_{gg} \left(-4I_1 + 4I_0 (\ln(1-x) - \ln(x) - 2) \right. \right. \\ & + 2\ln^2(1-x) + 2\ln^2(x) - 4\ln(x)\ln(1-x) - 8\ln(1-x) \\ & + \left. \frac{11}{3}\ln(x) + 2\frac{\pi^2}{6} + 4 \right) + \ln(x) \left(\frac{11}{3}x^2 + \frac{23}{6}x + \frac{23}{6} + \frac{11}{3x} \right) \\ & \left. - \frac{22}{3}x^2 + \frac{24}{3}x - \frac{25}{3} + \frac{22}{3x} \right]. \end{aligned}$$

Virtual contribution to NLO P_{gg} kernel in NPV

The virtual graph in New PV prescription:



$$\Gamma_{gg, NPV}^{(d_V)}(x, \epsilon) = C_S^{(d_V)} P_{gg} \left[\frac{1}{\epsilon^2} (1 + \epsilon \ln(1-x)) \tilde{Z}_{GS}^V \right. \\ \left. + \frac{1}{\epsilon} \left(4l_0 \ln(1-x) + 8l_0 \ln(x) - 16l_1 + 4 \ln^2(x) \right. \right. \\ \left. \left. + 12 \frac{\pi^2}{6} - \frac{134}{9} \right) \right] - C_S^{(d_V)} \frac{1}{3\epsilon} x$$

$$\tilde{Z}_{GS}^V = 12l_0 + 4 \ln(1-x) + 4 \ln(x) - \frac{22}{3},$$

Inclusive sum of Real and Virtual graphs $\tilde{\Gamma}_{gg}^{(d)}$
 identical as in standard PV scheme

**Both schemes, PV and New PV,
 give the same P_{qq} and P_{gg} kernels**

Summary



- ▶ We modified the way of using the PV prescription in the light-cone gauge by applying it to all the singularities in the plus component of the integration momentum.
- ▶ In the New PV prescription the NLO kernels P_{qq} and P_{gg} are reproduced correctly.
- ▶ Partial results of four graphs contributing to the kernels, differ in PV and New PV: the $1/\epsilon^3$ poles, present in PV, are replaced by $(1/\epsilon) \ln^2 \delta$ etc.
- ▶ Real graphs, free of $1/\epsilon^3$ and $1/\epsilon^2$ poles, can be calculated in four dimensions and are usable for the stochastic simulations.
- ▶ The higher order poles cancel separately for real and for virtual components.